

Precision Density Measurements Near the Helium Lambda Transition Using High-Q Microwave Cavities

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Precision Density Measurements Near the Helium Lambda Transition Using High-Q Microwave Cavities D. M.

STRAYER, Jet Propulsion Laboratory, Caltech, W. JIANG, N.-C. YEH, N. ASPLUND, Caltech — A new experimental approach for high-precision density measurements of liquid helium near the lambda transition is proposed. Using a high-Q Nb microwave cavity ($Q \sim 10^{10}$) and the high-resolution thermometry (HRT), the changes in the density of helium that fills the cavity can be detected to high precision by accurate measurements of the resonant frequency shift (Δf) as a function of the temperature. Since the frequency shift provides direct information for the changes in the dielectric constant, and since the dielectric constant is related to the density through the Clausius-Mossotti relation, the capability of high resolution frequency measurements (to one part in 10^{13}) will enable us to resolve density changes to one part in 10^{10} . Numerical calculations have been performed to demonstrate the feasibility of this approach for mapping out the density profile of liquid helium which couples to the TE modes of a microwave cavity. For temperatures very near the lambda transition, a superfluid-normal fluid interface develops inside the cavity. A numerical deconvolution technique is established to resolve the helium density profile in the cavity. Preliminary experimental data using a TM₀₁₀ niobium cavity and with microkelvin temperature resolutions will be presented.

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Prefer Oral Session
 Prefer Poster Session

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Motivation

- Test of the Critical Relation of β_P near the Lambda Transition

$$\begin{aligned}\beta_P &= A|t|^{-\alpha}(1 + D|t|^x + \dots) + B \\ &\equiv -\frac{1}{\rho} \frac{\partial \rho}{\partial T} \Big|_P\end{aligned}$$

High precision measurement of $\rho(T)$

$$\implies \beta_P(T) \implies \alpha, A, B, D, x$$

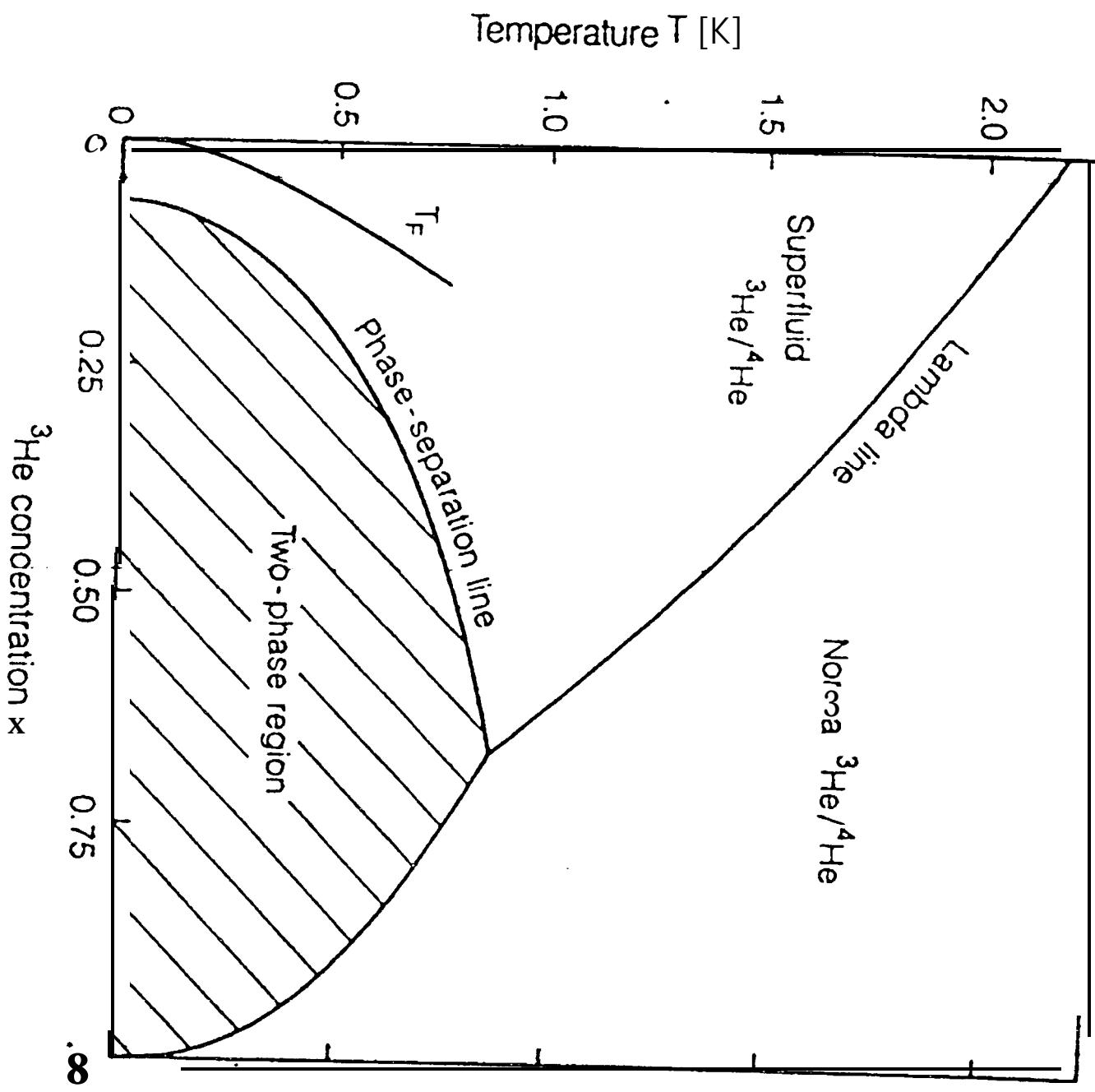
- Related Experimental Work:

| | Physical Quantity | Sensitivity ($\epsilon_\lambda \equiv 1 - (T/T_\lambda) $) |
|--|---|---|
| Van Degrift & Pellam 1974 | Capacitor (C) | $\delta C/C \sim 10^{-9}, \epsilon_\lambda \sim 10^{-5}$ |
| Muellel, Ahlers, & Pobell 1976 | thermal expansion coefficient (β_P) | $\delta P \sim 0^{-7}$ bar, $\epsilon_\lambda \sim 10^{-7}$ |
| Lips, <i>et al</i> (to be published in PRL) | heat capacity (C) | $\delta Q/Q \sim 10^{-4}, \epsilon_\lambda \sim 10^{-9}$ |
| this technique | resonant frequency shift ($\Delta f/f_0$) | $\delta f/f \sim 10^{-13}, \epsilon_\lambda \sim 10^{-9}$ |

- Other Applications:

$^3He - ^4He$ mixture: $\rho_{mix} = x\rho_{^3He} + (1-x)\rho_{^4He}$

Precision measurements of $\rho \implies$ precision measurement of x .



Technical Approach

- Clausius-Mossotti Relation ($\epsilon \iff \rho$)

$$\frac{\epsilon - 1}{\epsilon + 2} = \frac{4\pi\alpha_0}{3M}\rho$$

ϵ : dielectric constant

α_0 : polarizability

(assumed to be constant near the lambda transition])

M : molecular weight

ρ : density

- In a microwave cavity, the electric fields couple to the dielectric constants ϵ

- Cavity resonant frequency shift(Δf) due to the small change of ϵ of liquid helium:

$$\frac{\Delta f}{f_0} = \frac{f - f_0}{f_0} = \frac{\int_{V_0} (\epsilon - \epsilon_0) |E_0|^2 dV}{\int_{V_0} \epsilon_0 |E_0|^2 dV}$$

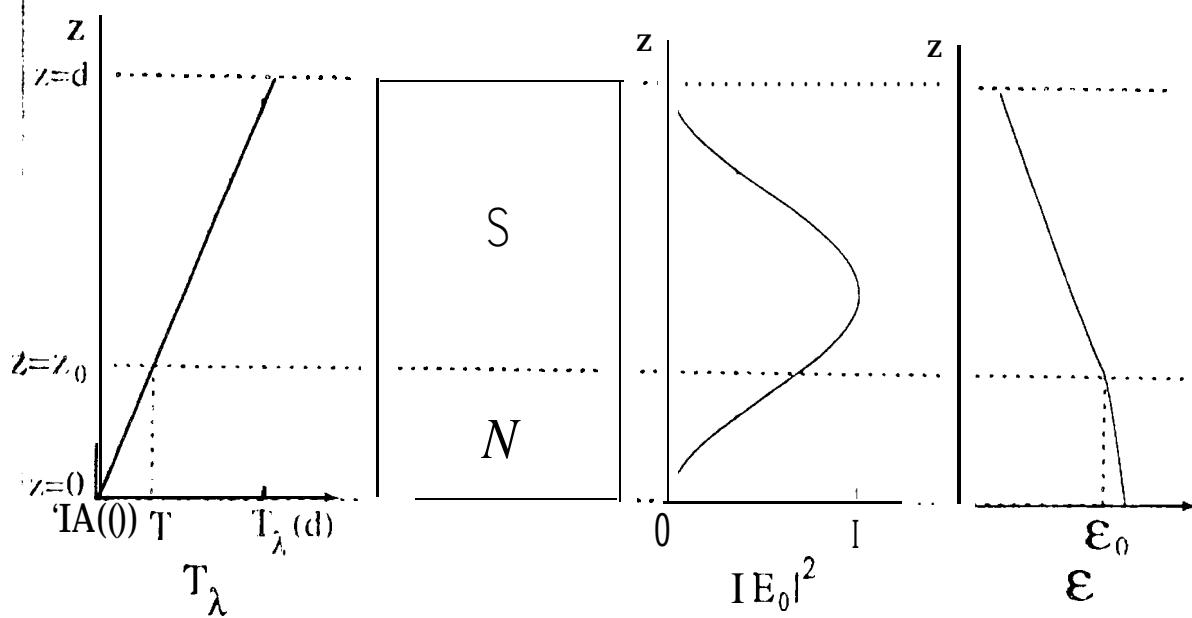
f_0 : resonant frequency of the resonator
at a reference temperature

V_0 : volume of the cavity

E_0 : the electric field of the resonant mode

ϵ_0 : dielectric constant(s) at the reference temperature

● Gravity [Effect]



Consider a cylindrical cavity maintained at a uniform temperature T ,

$$T_{\lambda}(z) = T(0) - \frac{1}{2} \gamma_0 z$$

$$\gamma_0 = 1.273 \times 10^{-6} \text{ K/cm}$$

Ref: Ahlers, Phys. Rev. 171, 275 (1968).

If $T_{\lambda}(0) < T < T_{\lambda}(d)$, a superfluid/normal fluid interface appears inside the cavity at the position

$$z_0 = \frac{T - T_{\lambda}(0)}{\gamma_0}$$

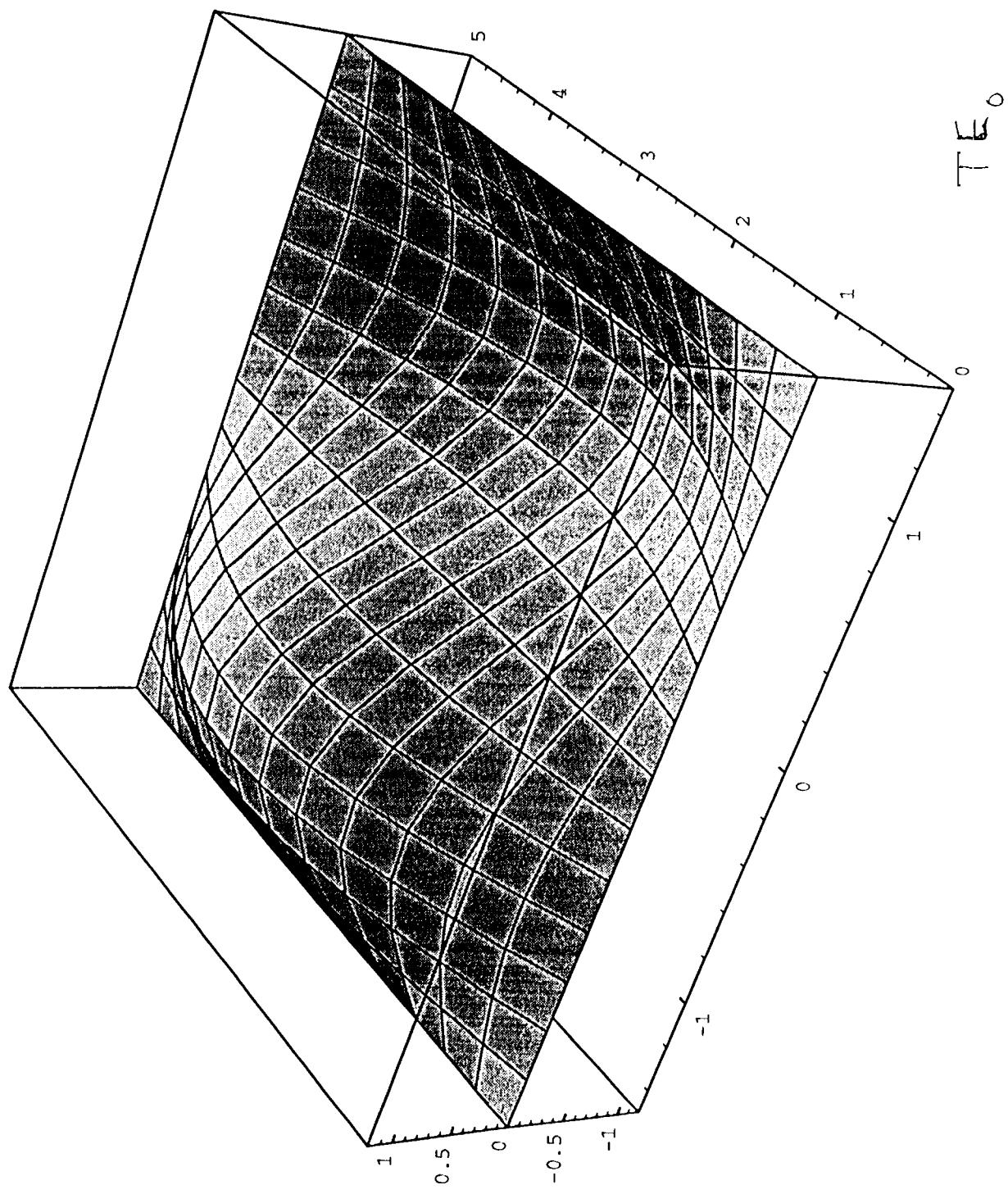
• Deconvolution for Resolving $\rho(z, T)$ near the Lambda Transition

$$\frac{\Delta f}{f_0}(T) \xrightarrow{\text{deconvolution}} \varepsilon(z, T) \xrightarrow{\text{Clausius-Mossotti relation}} \rho(z, T)$$

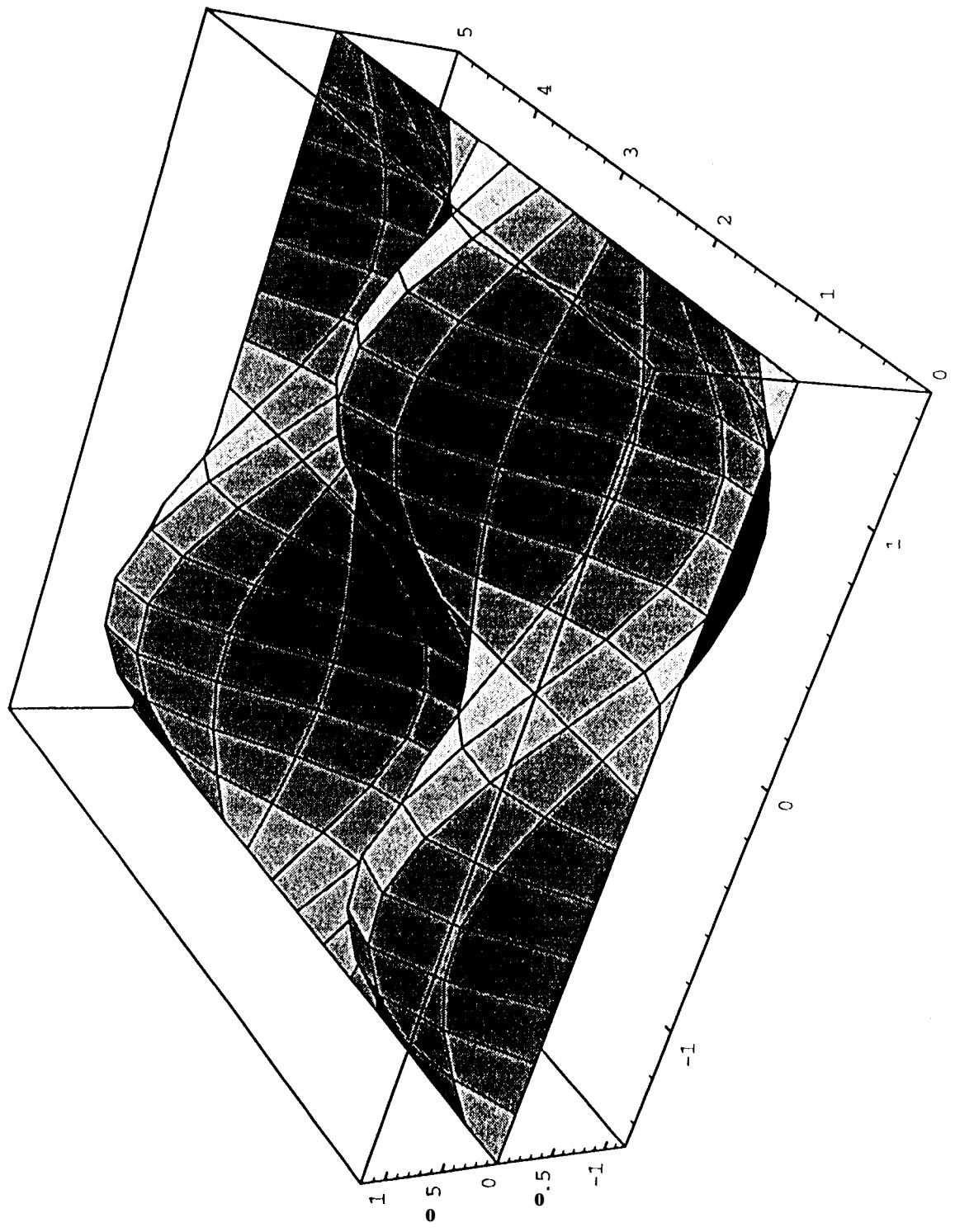
For $TE_{nm\ell}$ mocks,

$$\vec{E} = E_\theta \hat{e}_\theta, E_\theta = \frac{ik\eta a H_0}{P'_{nm}} \left(\frac{\bar{P}'_{nm} r}{a} \right) \sin \left(\frac{\ell \pi z}{d} \right)$$

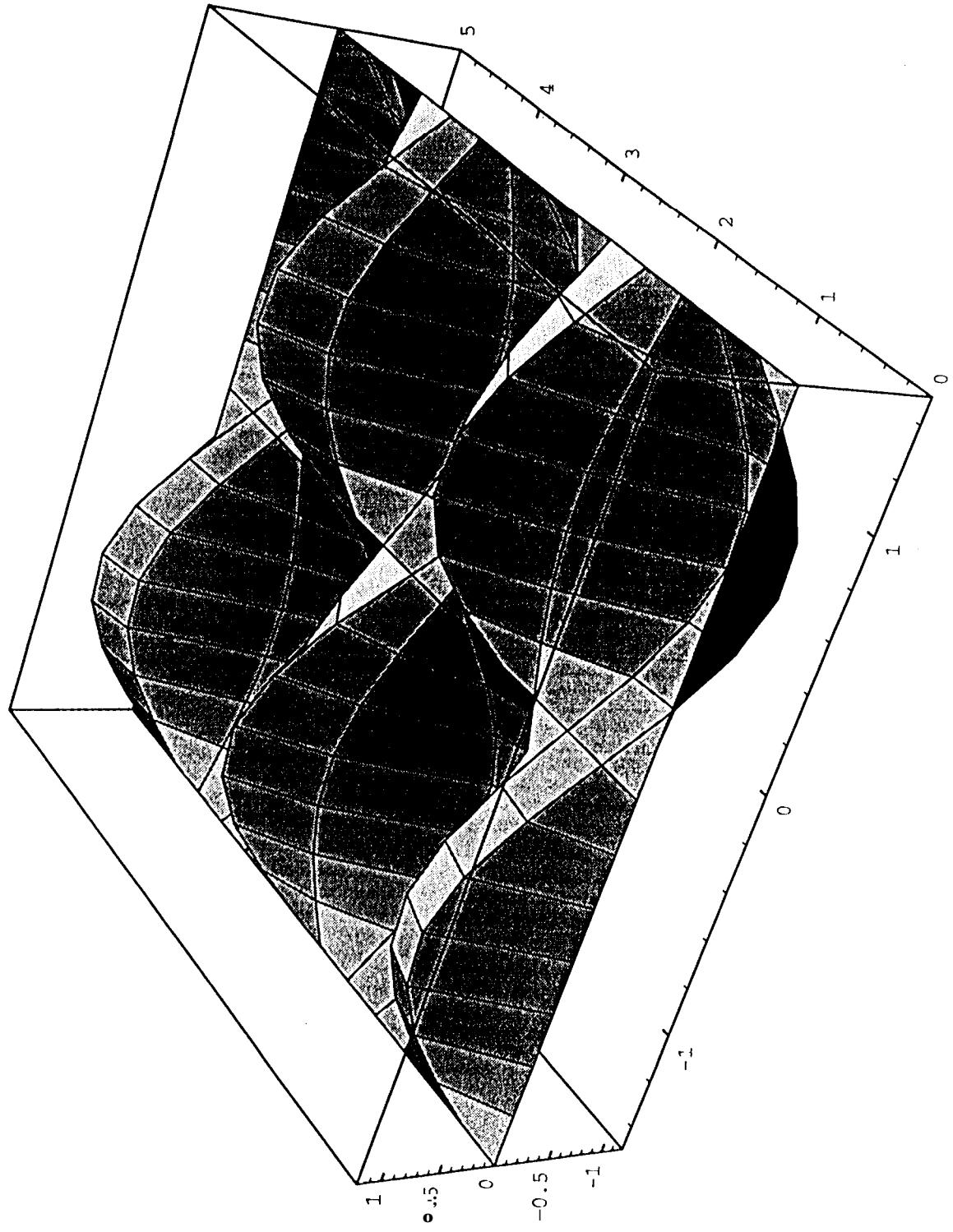
$$\frac{\Delta f}{f_0}(T) = \frac{f - f_0}{f_0} = \frac{\int_0^d [\varepsilon(z, T) - \varepsilon(z, T_\lambda(0))] \sin^2 \left(\frac{\ell \pi z}{d} \right) dz}{\int_0^d \varepsilon(z, T_\lambda(0)) \sin^2 \left(\frac{\ell \pi z}{d} \right) dz}$$



$\overline{TE}_{0(3)}$



\overline{TE}_0 5



Experimental

—— High-Q Nb microwave cavity:

$Q \sim 10^9$ near 2.2 K

—— High-Resolution Thermometry (HRT):

SQUID magnetic susceptibility:

Ref: Lipa et al, *Physica B+C* 107, 331(1981)

Temperature resolution: $\sim 10^{-10}$ K

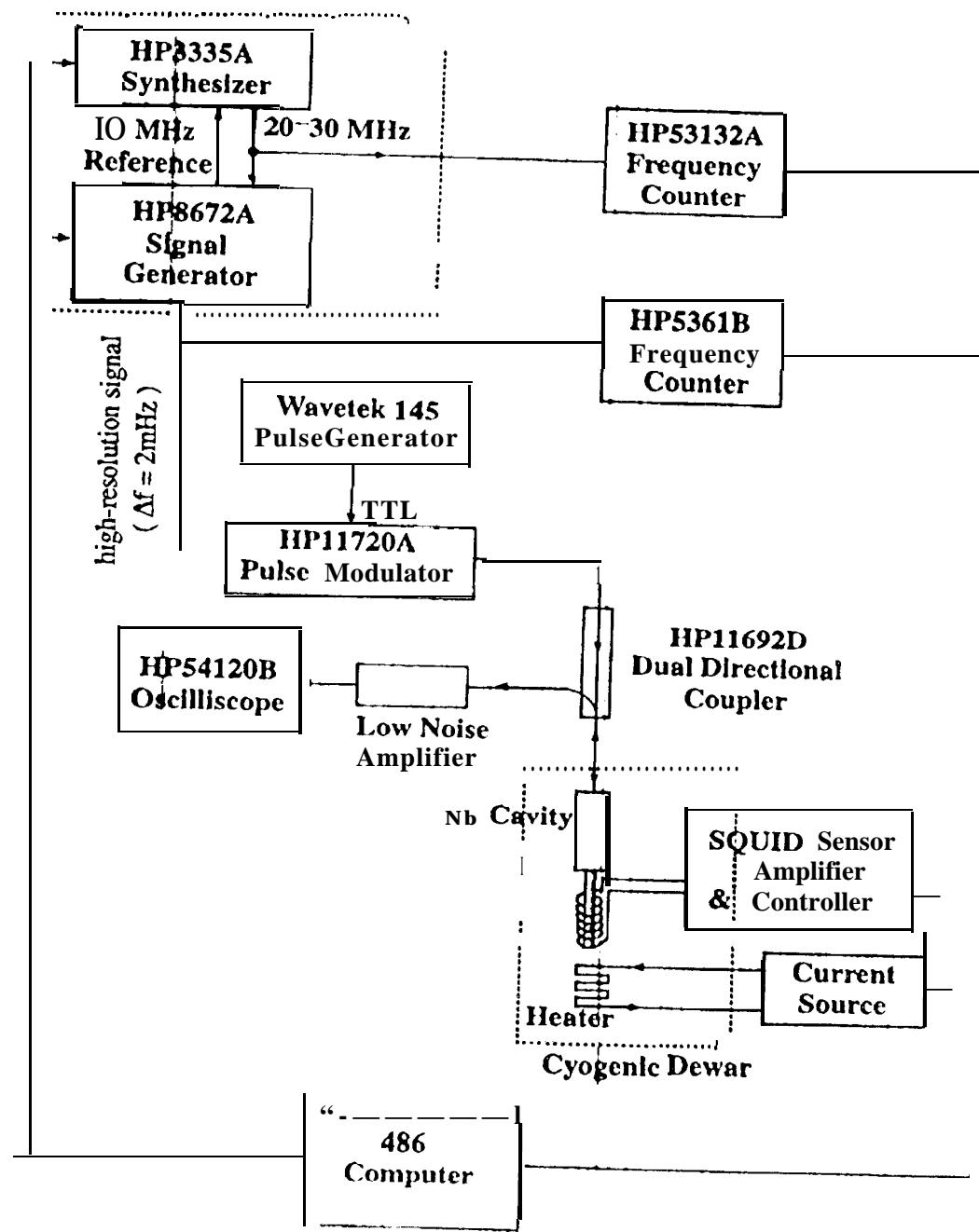
Temperature stability: $\sim 10^{-9}$ K

—— High-resolution frequency source:

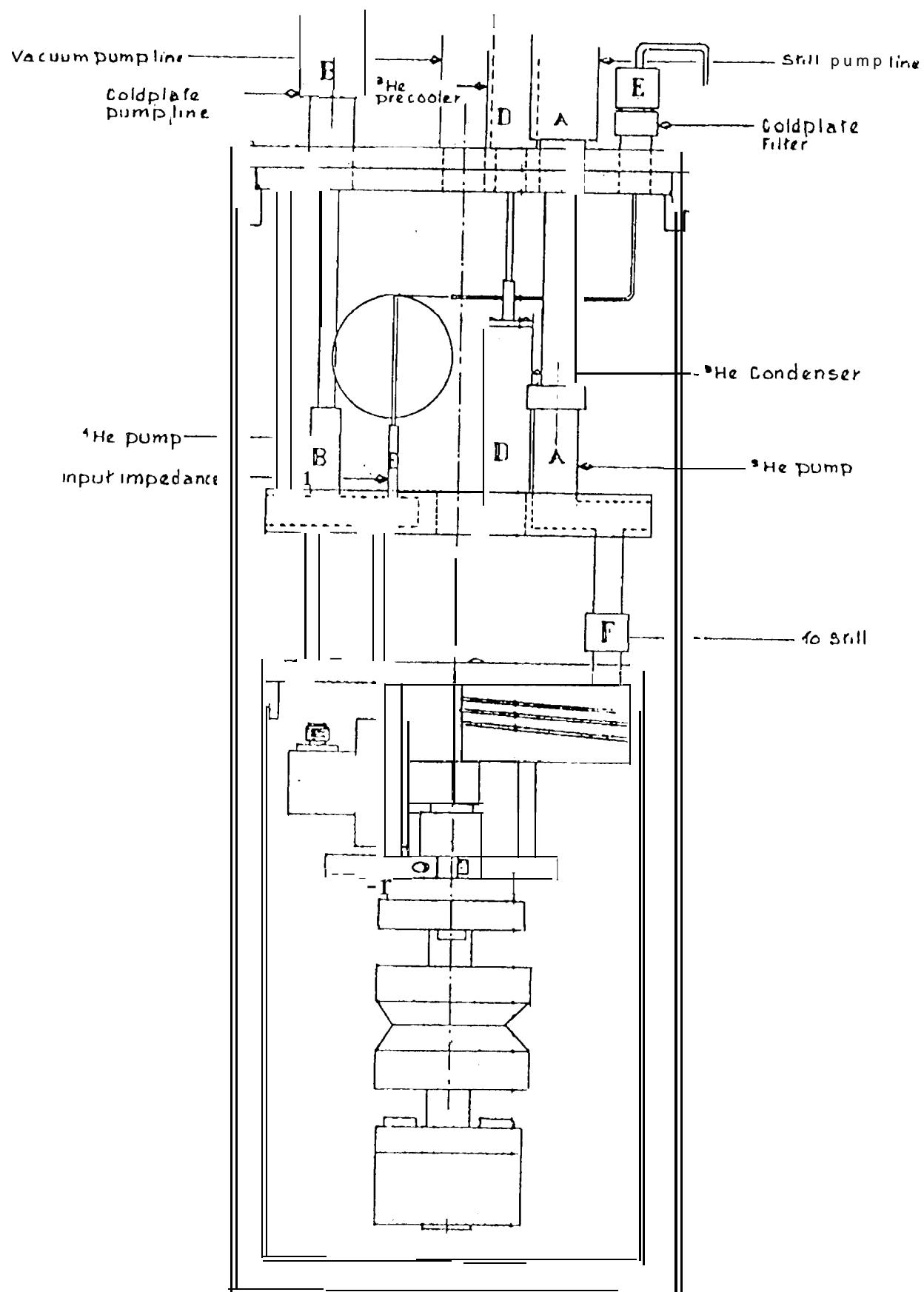
HP high-stability source: $\delta f/f \sim 10^{-13}$

phase-lock loop technique: $\delta f/f \sim 10^{-15}$

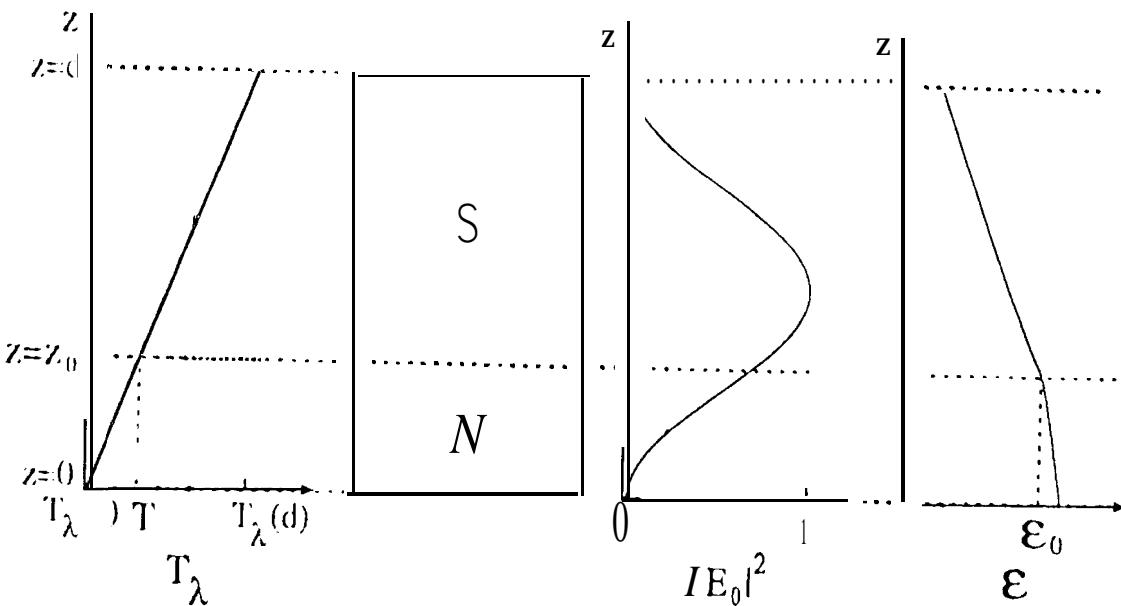
• Block Diagram of Resonant Frequency Measurement Using a High-Q Microwave Cavity



Schematic Measurement Cell Design



Deconvolution Algorithm



For $T_\lambda(0) < T < T_\lambda(z)$,

a given temperature $T \implies$ an interface position z_0

Choose

$$T_j = T_\lambda(0) + \gamma_0 z_j, j = 1, \dots, N,$$

such that $z_j = j \frac{d}{N}$

Discretization

$$\frac{\Delta f}{f_0}(T) = \frac{f_0 - f(0)}{f(0)} = -\frac{\int_0^d [\varepsilon(z, T) - \varepsilon(z, T_\lambda(0))] \sin^2(\frac{\ell\pi z}{d}) dz}{\int_0^d \varepsilon(z, T_\lambda(0)) \sin^2(\frac{\ell\pi z}{d}) dz}$$

$$\begin{aligned} &\Leftrightarrow \left[1 - \frac{\Delta f}{f_0}(T_j) \right] \sum_{i=1}^N [\varepsilon(t_i^0) - \varepsilon_0] \sin^2\left(\frac{\ell\pi z_i}{d}\right) - \sum_{i=1}^N [\varepsilon(t_i^j) - \varepsilon_0] \sin^2\left(\frac{\ell\pi z_i}{d}\right) \\ &= \frac{\Delta f}{f_0} \frac{\varepsilon_0 N}{2} \end{aligned}$$

where

$$\varepsilon_0 \equiv \varepsilon(T = T_\lambda),$$

$$t_i^j = t(T_j, z_i) \equiv T_j - T_\lambda(z_i) = (j - i) \frac{\gamma_0 d}{N}$$

$$z_i = i \frac{d}{N}$$

$$i = 1, \dots, N, j = 1, \dots, N$$

\iff Linear algebra problem

$$\mathbf{A} \mathbf{x} = \mathbf{b} \implies \mathbf{x} = (\mathbf{A}^T \mathbf{A})^{-1} \mathbf{A}^T \mathbf{b}$$

where

$$\mathbf{x} \equiv \varepsilon - \varepsilon_0, \quad \mathbf{b} = \frac{\varepsilon_0 N}{2} \mathbf{C}$$

$$\varepsilon = \begin{pmatrix} \varepsilon(-dt) \\ \varepsilon\left(-\frac{N-1}{N}dt\right) \\ \vdots \\ \vdots \\ \varepsilon(dt) \end{pmatrix}, \quad \mathbf{C} = \begin{pmatrix} \frac{\Delta f}{f_0}(T_1) \\ \frac{\Delta f}{f_0}(T_2) \\ \vdots \\ \vdots \\ \frac{\Delta f}{f_N}(T_N) \end{pmatrix}$$

$$\mathbf{A} = \begin{pmatrix} (1 - C_1)\mathbf{u}^T & 0 & \dots & 0 \\ (1 - C_2)\mathbf{u}^T & 0 & \dots & 0 \\ \vdots & 0 & \dots & 0 \\ \vdots & 0 & \dots & 0 \\ (1 - C_N)\mathbf{u}^T & 0 & \dots & 0 \end{pmatrix} - \begin{pmatrix} 0 & \mathbf{u}^T & 0 & \dots & 0 \\ 0 & 0 & \mathbf{u}^T & \dots & 0 \\ \vdots & \dots & \dots & \dots & 0 \\ \vdots & \dots & \dots & \dots & 0 \\ 0 & 0 & \dots & \dots & \mathbf{u}^T \end{pmatrix}, \quad \mathbf{u} = \begin{pmatrix} \sin^2\left(\frac{\ell\pi z_1}{d}\right) \\ \sin^2\left(\frac{\ell\pi z_2}{d}\right) \\ \vdots \\ \vdots \\ \sin^2\left(\frac{\ell\pi z_N}{d}\right) \end{pmatrix}$$

With $z_i \neq 1$ & $3, 2N$ unknowns: $\varepsilon - \varepsilon_0$ (assume ε_0 is known), and $2N$ equations.

• Computer Simulation

$$\rho(t) = \begin{cases} 0.146081 - 0.1451189 \times 10^{-6.0} \times (2630.9 t \\ - 7768.2 t \log_{10}|t| + 20547.1 t^2 + 15963.2 t \log_{10}|t| \\ + 24525.3 t^3 - 12728 t^2 t \log_{10}|t|), & \text{for } t < 0; \\ 0.146081, & t = 0; \\ 0.146081 + 0.1451189 \times 10^{-6.0} \times (-31288.6 t \\ - 8086.1 t \log_{10}t - 10738.1 t^2 - 3010.4 t^3), & \text{for } t > 0. \end{cases}$$

where $t \equiv T - T_\lambda(z)$

$$\alpha_0^* = 0.123363 N_A = 0.204849 \text{ cm}^3/\text{molecule}$$

$$M^* = 6.6424 \times 10^{-24} \text{ g}$$

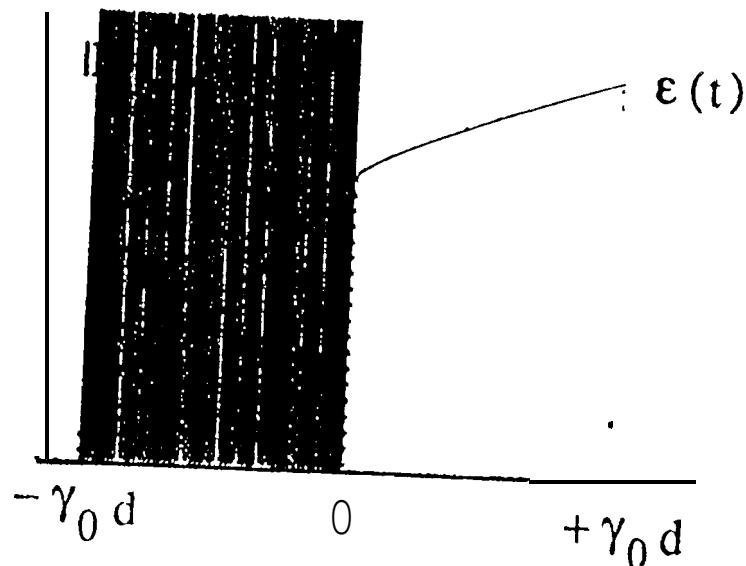
$$\Rightarrow \varepsilon(t) = \frac{1 + 2a\rho}{1 - a\rho}, \quad a \equiv \frac{4\pi\alpha_0}{3M}$$

$$\Rightarrow \frac{\Delta f}{f_0}(T) = \frac{f - f_0}{f_0} = -\frac{\int_0^d [\varepsilon(z, T) - \varepsilon(z, T_\lambda(0))] \sin^2(\frac{\ell\pi z}{d}) dz}{\int_0^d \varepsilon(z, T_\lambda(0)) \sin^2(\frac{\ell\pi z}{d}) dz}$$

* Ref: Donnelly, Riegelmann & Barenghi,
 "The Dielectric Properties of Liquid Helium
 at the Saturated Vapor Pressure", 1953

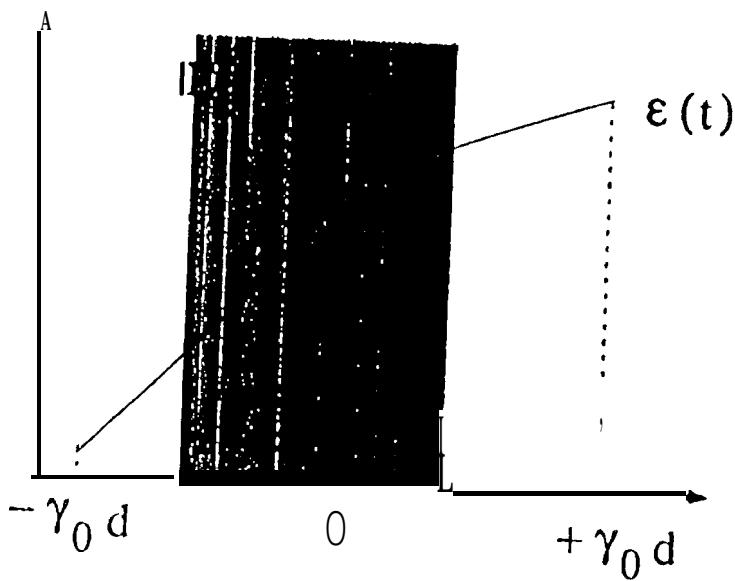
1
1
1
1
1

$T = T_\lambda(0)$



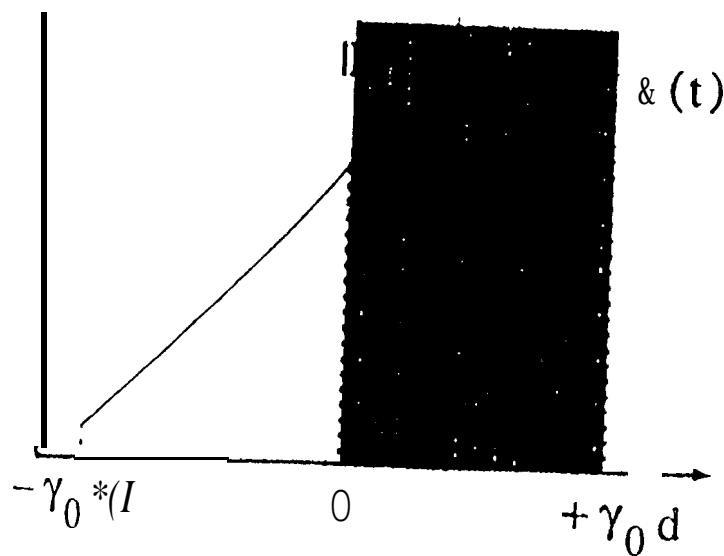
$t = T - T_\lambda(z)$

$T_\lambda(0) < T < T_\lambda(d)$

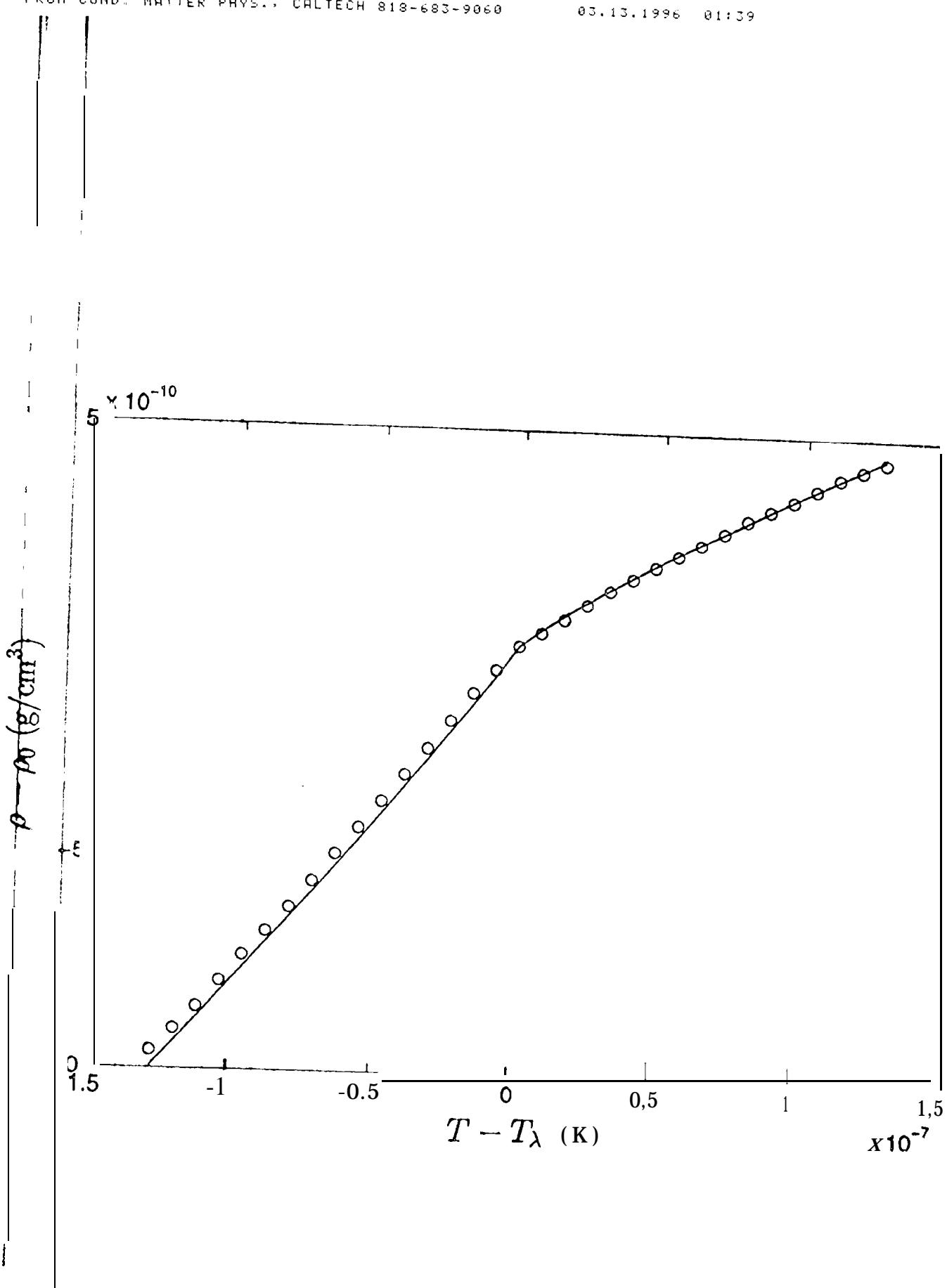


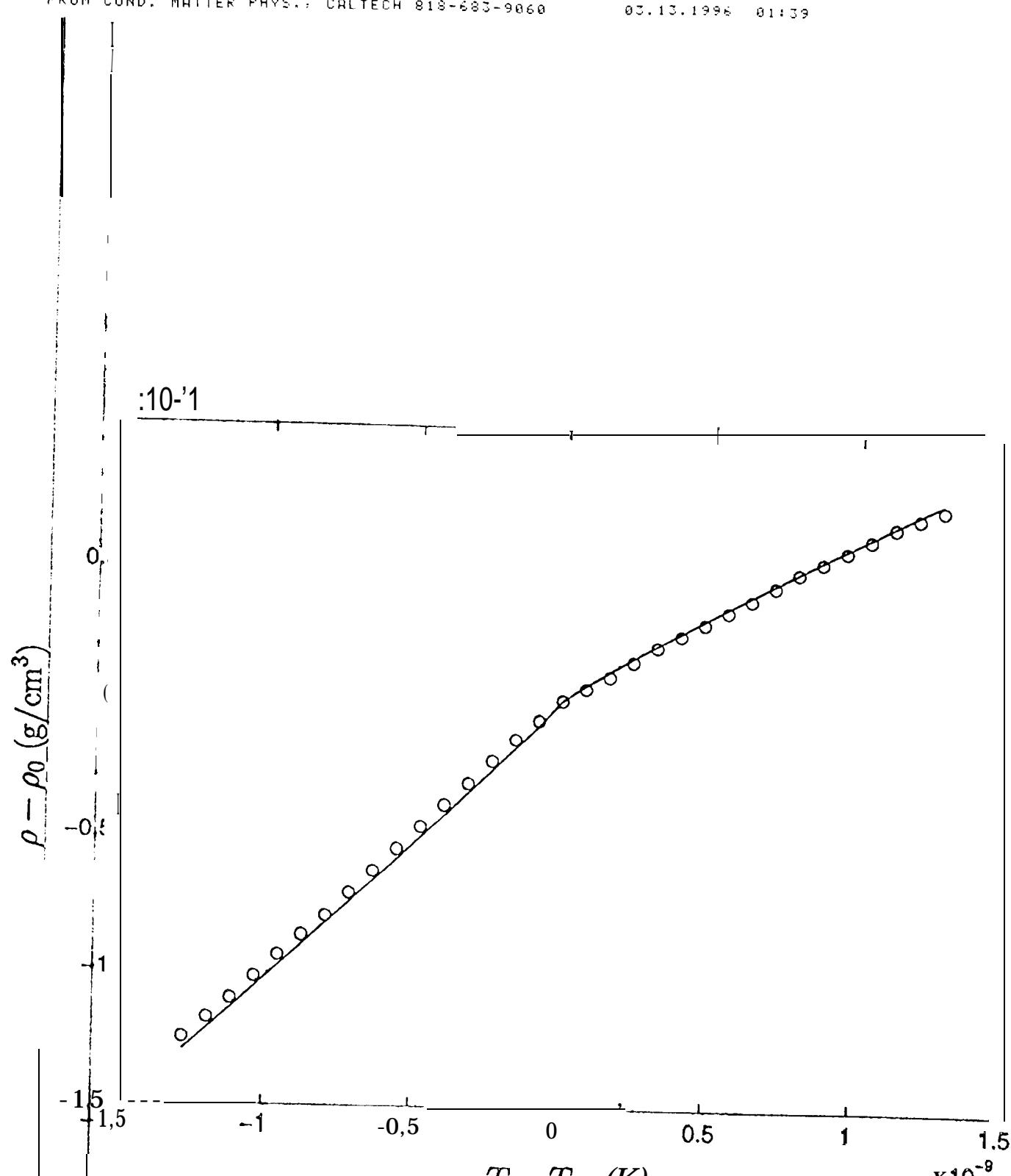
$t = T - T_\lambda(z)$

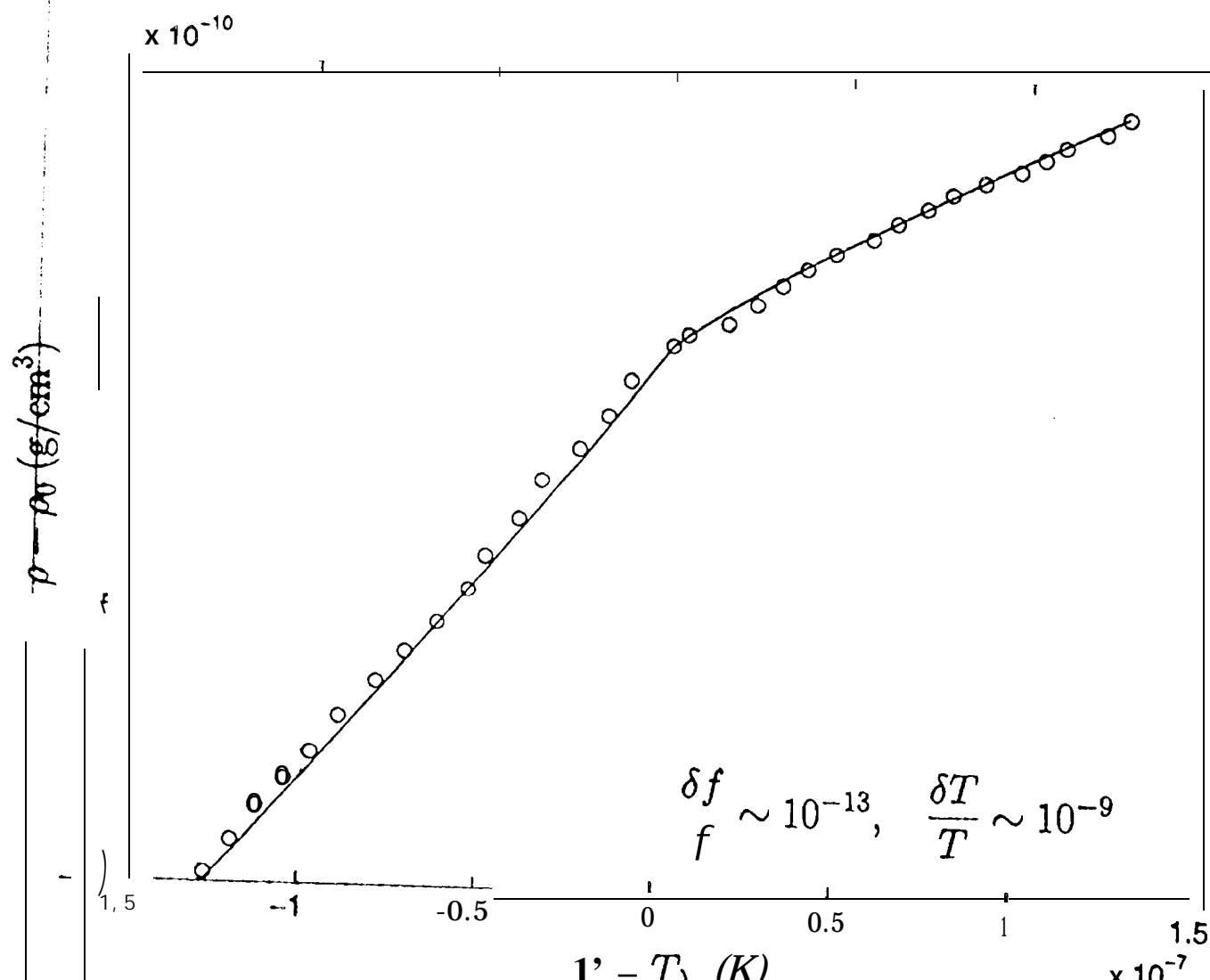
$T \neq T_\lambda(d)$

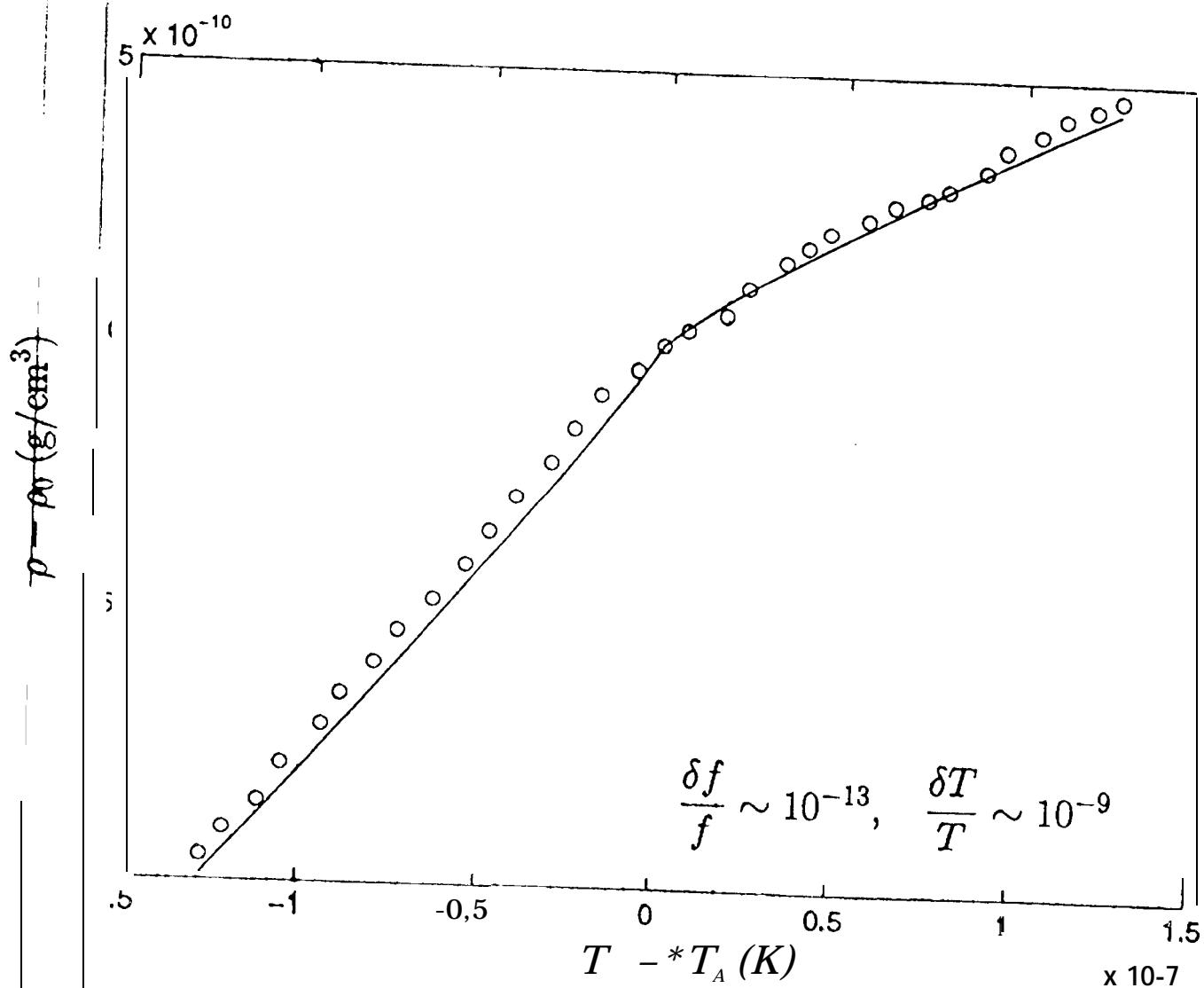


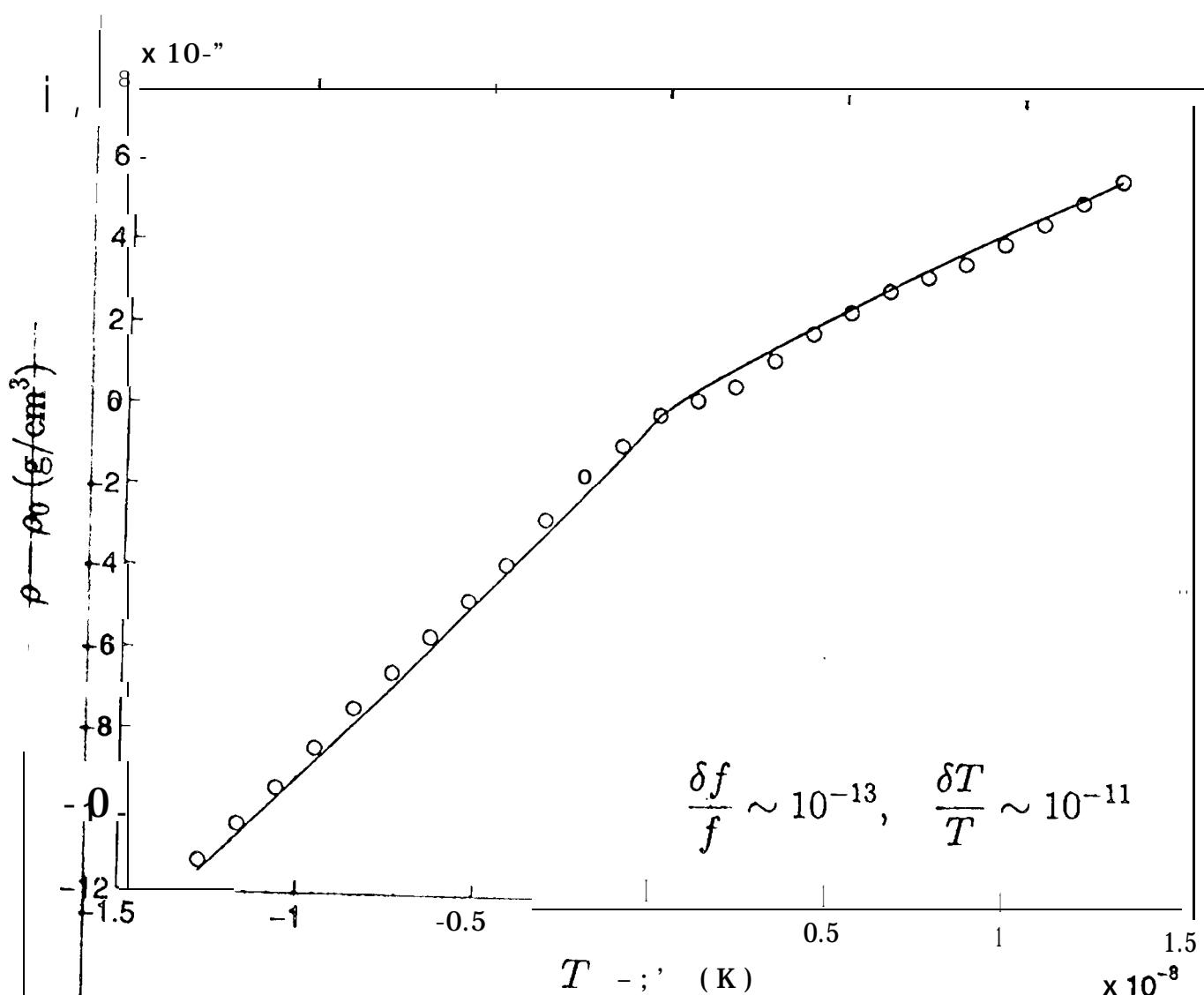
$t = T - T_\lambda(z)$

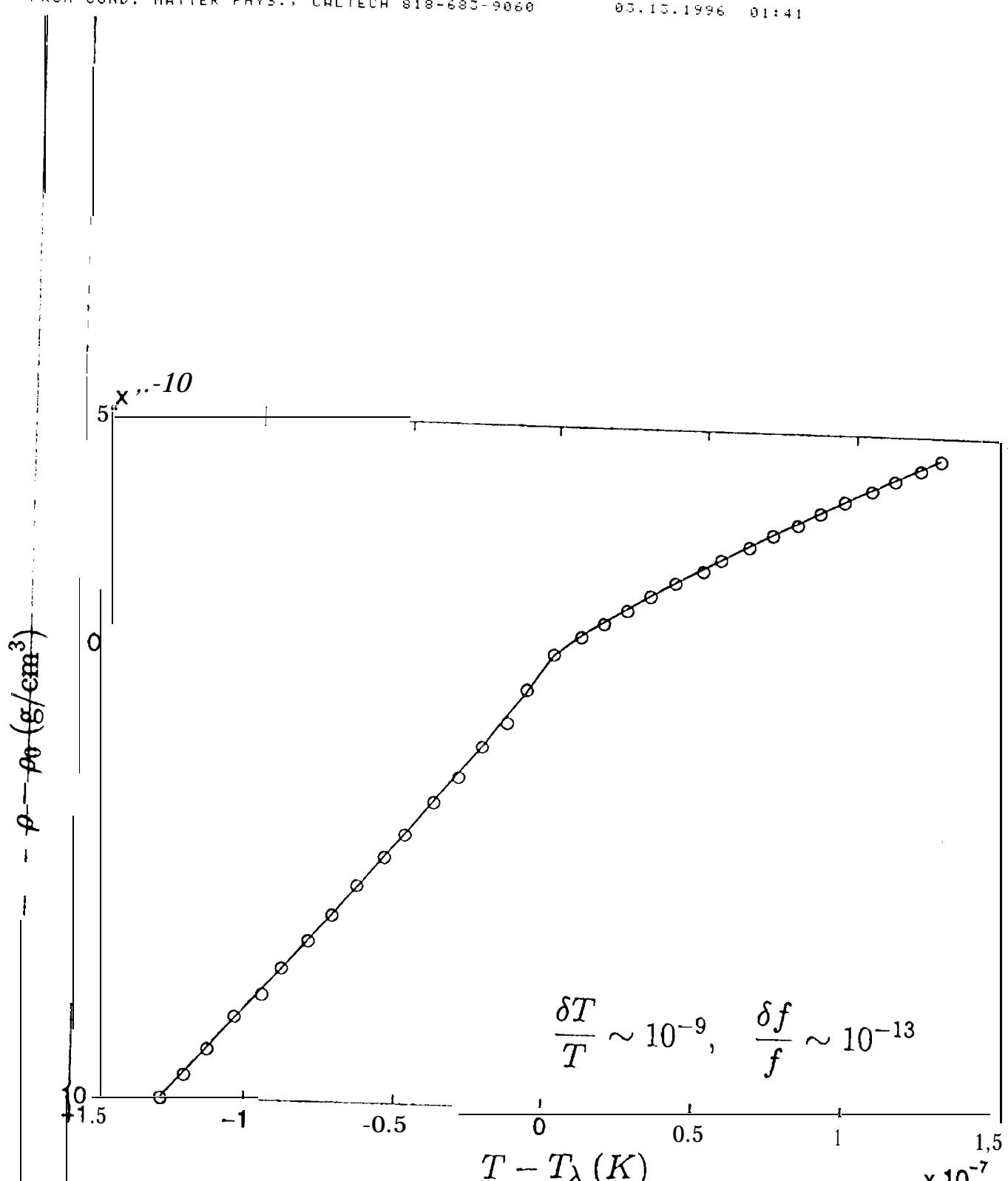




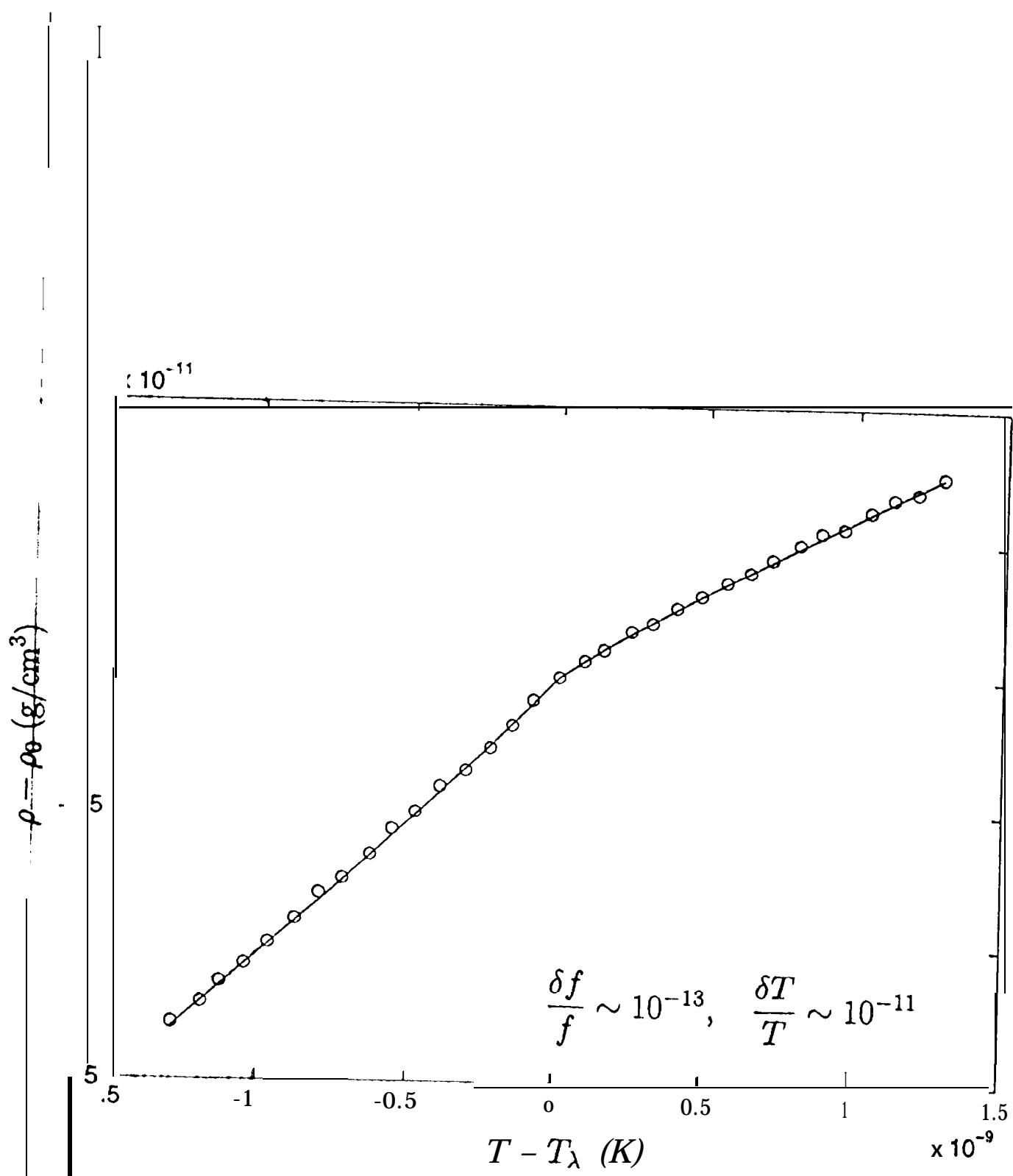








microgravity environment



Microgravity
Environment

- Expected Resolution in Density Measurement

| | $\delta f/f$ | $\delta T/T$ | $6p/p$ |
|--------------------------|------------------------------------|-----------------------------------|--|
| earth-bound laboratory | $\sim 10^{-13}$ $\sim 10^{-13}$ | $\sim 10^{-9}$ $\sim 10^{-11}$ | $\sim < 3.5 \times 10^{-10}$ $\sim < 0.8 \times 10^{-10}$ |
| microgravity environment | -10^{-13} -10^{-13} | $\sim 10^{-9}$ -10^{-11} | $\sim < 2.5 \times 10^{-10}$ $\sim < 2.5 \times 10^{-12}$ |

Summary

- High precision measurements of helium density near the lambda transition

High-Q Nb cavity ($Q \sim 10^9$ near 2.2 K)

High-resolution thermometry ($|T - T_\lambda| \sim 10^9$ K)

High-resolution frequency source ($\delta f/f \sim 10^{-13}$)

$$\Rightarrow \delta\rho/\rho \sim 10''^{10}$$

- Deconvolution algorithm

$$\frac{\Delta f}{f_0}(T) \Rightarrow \epsilon(T, z) [\Rightarrow \rho(T, z)]$$

- * Potential of the Technique:

Higher precision density measurements in the microgravity environment: $\delta T/T \sim 10^{-11} \Rightarrow \delta p/p \sim 10^{-12}$